

# Self-testing Quantum Supermaps

## with an application to the Quantum Switch

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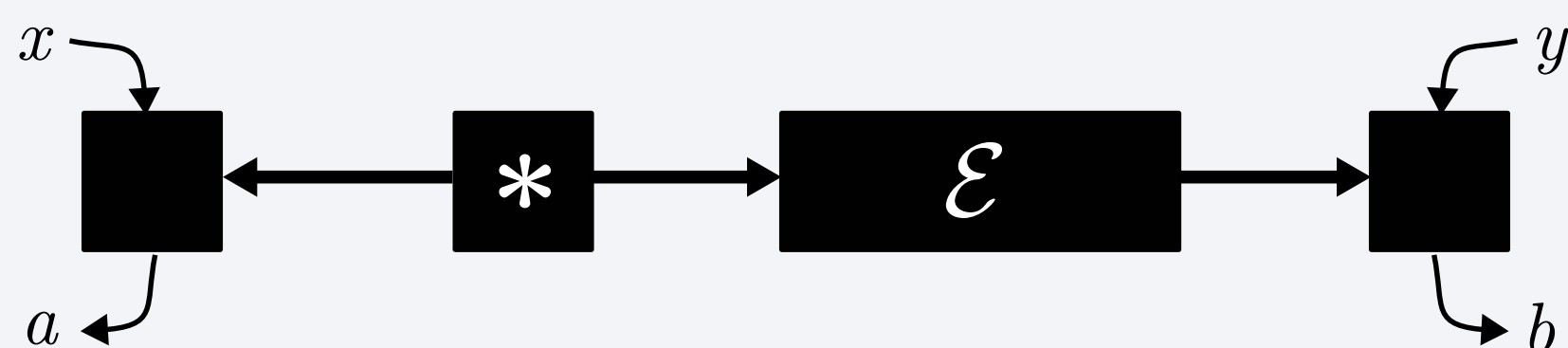
### Motivation

- Ongoing debate regarding experiments implementing the Quantum Switch and the ability to certify its causal indefiniteness in a device-independent setting
- Use self-testing techniques to certify a Quantum Switch in a black-box way
- Framework of self-testing needs to be generalised to quantum supermaps

### Self-testing

- Determine the most accurate physical description of a device possible without prior knowledge of the internal workings of the apparatuses involved
- **Black-box setting**: devices used are uncharacterised but their connections are trusted
- Characterise device from the **observed correlations** and the **network structure**

**Example:** self-testing a quantum channel [1]



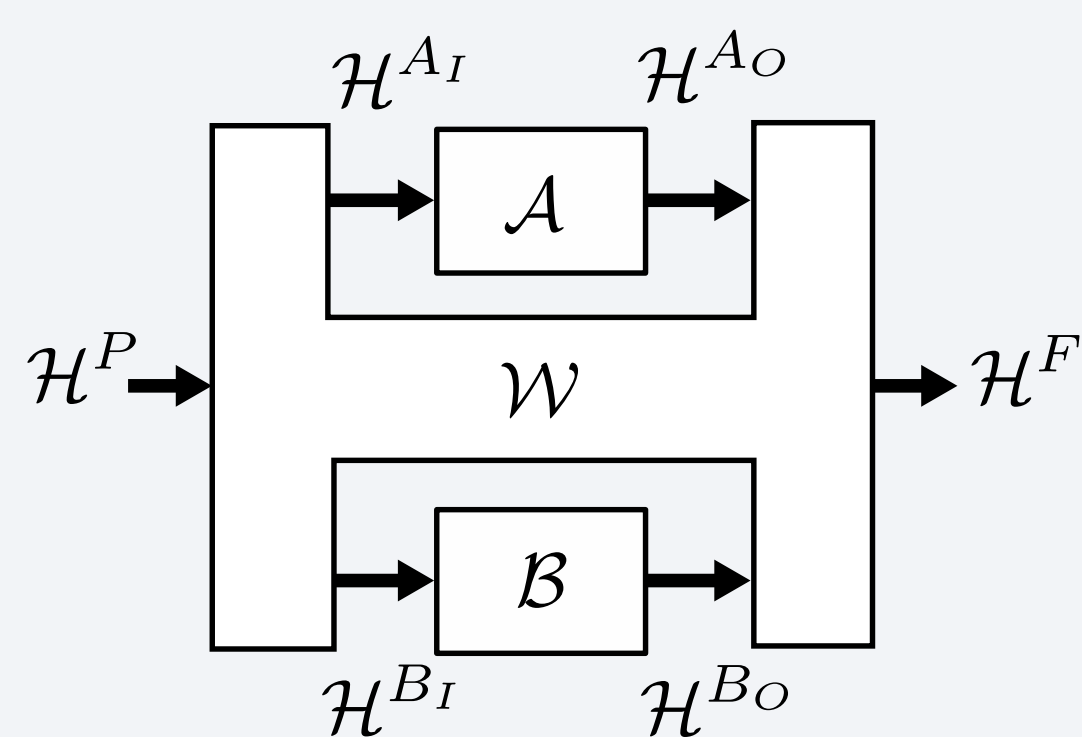
- Can only self-test up to DOFs undetectable in the black-box setting
  - Local change of basis, ancillas, etc. undetectable through correlations alone
- Desired correlations  $p(a, b|x, y) \implies \exists \Lambda_I, \Lambda_O$  such that:

$$\mathcal{H}^I \xrightarrow{\Lambda_I} \mathcal{E}_{\text{physical}} \xrightarrow{\Lambda_O} \mathcal{H}^O = \mathcal{H}^I \xrightarrow{\tilde{\mathcal{E}}_{\text{ideal}}} \mathcal{H}^O$$

### Quantum supermaps

- **Higher-order quantum operation**: transforms quantum maps into quantum maps

$$(\mathcal{A}, \mathcal{B}) \mapsto \mathcal{W}(\mathcal{A}, \mathcal{B})$$



- Represented in Choi picture as a **process matrix**:
 
$$W \in \mathcal{L}(\mathcal{H}^P \otimes \mathcal{H}^{A_I} \otimes \mathcal{H}^{A_O} \otimes \mathcal{H}^{B_I} \otimes \mathcal{H}^{B_O} \otimes \mathcal{H}^F)$$

- Choi matrices

$$A \in \mathcal{L}(\mathcal{H}^{A_I} \otimes \mathcal{H}^{A_O}) \text{ and } B \in \mathcal{L}(\mathcal{H}^{B_I} \otimes \mathcal{H}^{B_O})$$

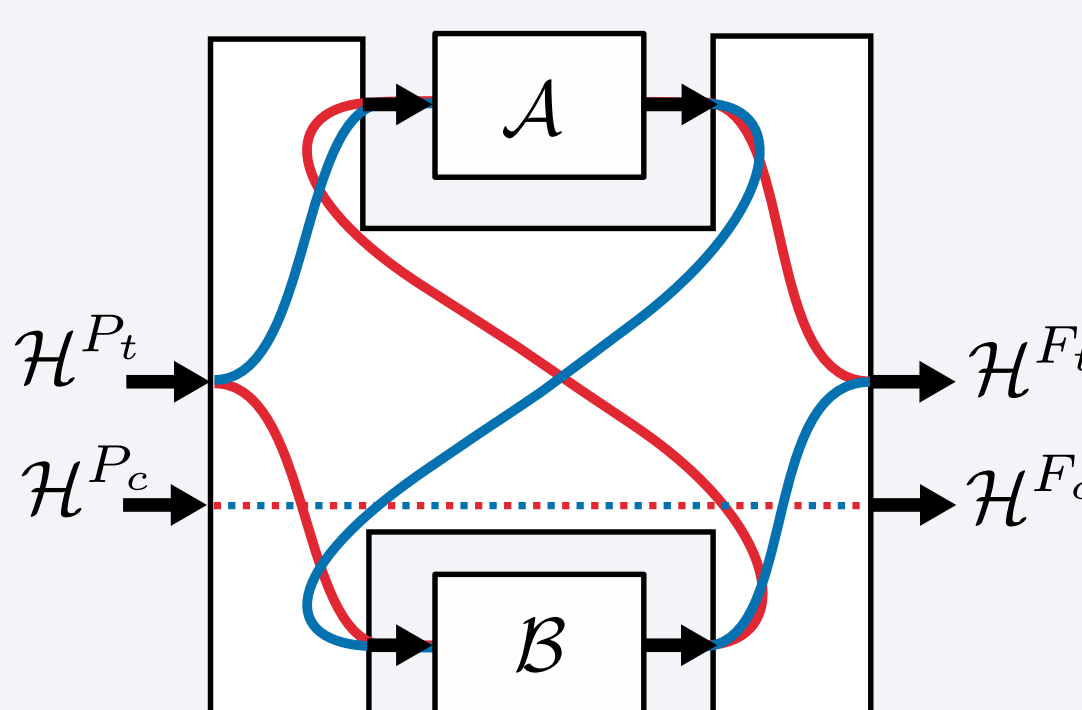
of  $\mathcal{A}$  and  $\mathcal{B}$  combined into

$$W * (A \otimes B) \in \mathcal{L}(\mathcal{H}^P \otimes \mathcal{H}^F)$$

### The Quantum Switch

- Some quantum supermaps describe causally indefinite compositions of  $\mathcal{A}$  and  $\mathcal{B}$ 
  - **Causal nonseparability**:  $\mathcal{W} \neq q\mathcal{W}^{A \prec B} + (1-q)\mathcal{W}^{B \prec A}$
- This property can be tested in both device-dependent (DD) and, sometimes, device-independent (DI) settings

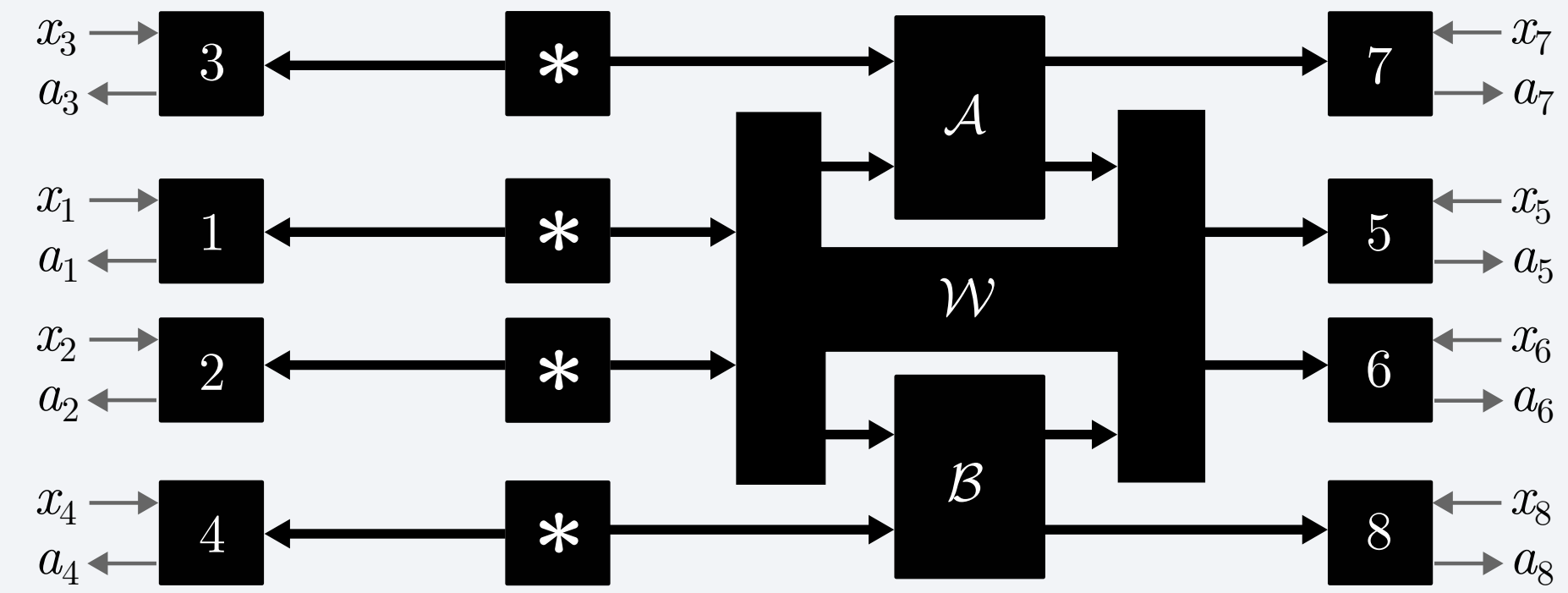
**Example:** The **Quantum Switch**



- Coherent quantum control of  $\mathcal{B} \circ \mathcal{A}$  and  $\mathcal{A} \circ \mathcal{B}$
- $\mathcal{W}_{\text{QS}} = |w_{\text{QS}}\rangle\langle w_{\text{QS}}|$  with
 
$$|w_{\text{QS}}\rangle = |0\rangle^{P_c} |\mathbb{1}\rangle^{P_t A_I} |\mathbb{1}\rangle^{A_O B_I} |\mathbb{1}\rangle^{B_O F_t} |0\rangle^{F_c} + |1\rangle^{P_c} |\mathbb{1}\rangle^{P_t B_I} |\mathbb{1}\rangle^{B_O A_I} |\mathbb{1}\rangle^{A_O F_t} |1\rangle^{F_c},$$
 where  $|\mathbb{1}\rangle = \sum_i |i, i\rangle \propto |\phi^+\rangle$
- The causal nonseparability of the Quantum Switch can be witnessed in a network DI setting [2, 3]
- But existing DI tests don't self-test the Quantum Switch

### Self-testing scheme

- Both the supermap and the devices plugged into it must be treated as black boxes



**Reference scenario:**

- Sources are maximally entangled qubit states  $|\phi^+\rangle$
- $\mathcal{A}$  and  $\mathcal{B}$  are SWAP gates
- The Quantum Switch then acts as a controlled swap on (half of) four  $|\phi^+\rangle$  states:

$$|\bar{\psi}_{\text{out}}\rangle = \frac{1}{4} \sum_{i,j,k=0}^1 \left( |00\rangle^{26} |ijk\rangle^{134} |kij\rangle^{578} + |11\rangle^{26} |ijk\rangle^{134} |jki\rangle^{578} \right)$$

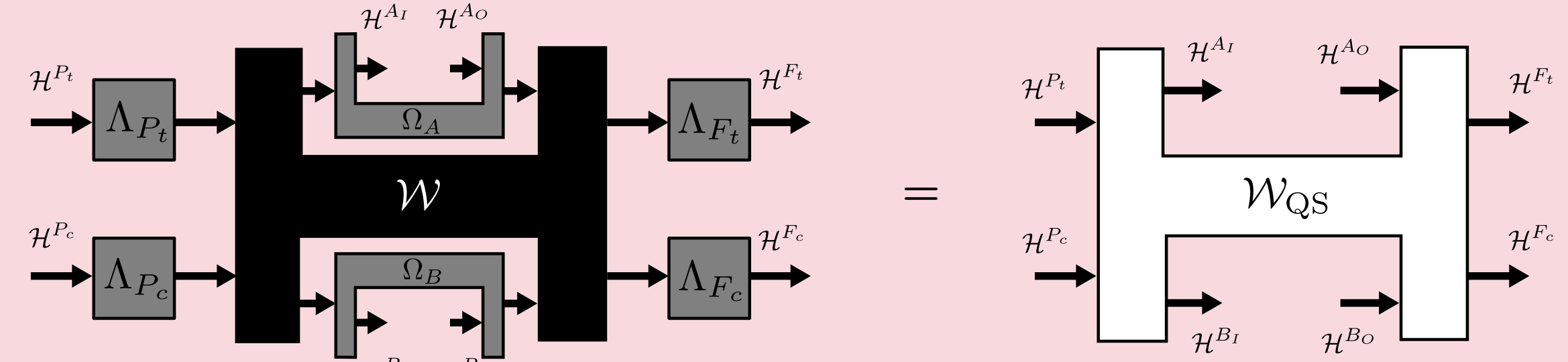
- **Idea:** self-test  $|\bar{\psi}_{\text{out}}\rangle$  with multiple permuted CHSH tests
  - Parties 1,2,3,4 measure  $M_0 = \sigma_z$ ,  $M_1 = \sigma_x$
  - Parties 5,7,8 measure  $M_0 = \frac{\sigma_z + \sigma_x}{\sqrt{2}}$ ,  $M_1 = \frac{\sigma_z - \sigma_x}{\sqrt{2}}$
  - Party 6 measures  $M_0 = \sigma_z$ ,  $M_1 = \sigma_x$ ,  $M_2 = \frac{\sigma_z + \sigma_x}{\sqrt{2}}$

**Theorem: Self-testing the input and output states**

- Given these statistics,  $\exists \Lambda_i : \mathcal{L}(\mathcal{H}_i) \rightarrow \mathbb{C}^2$  such that  $\otimes_{i=1}^8 \Lambda_i(\rho_{\text{out}}) = |\bar{\psi}_{\text{out}}\rangle\langle\bar{\psi}_{\text{out}}|$
- The same measurements self test  $\rho_{\text{in}}$ :  $\exists \Lambda'_i$  such that  $\otimes_{i=1}^8 \Lambda'_i(\rho_{\text{in}}) = |\phi^+\rangle\langle\phi^+|^{\otimes 4}$

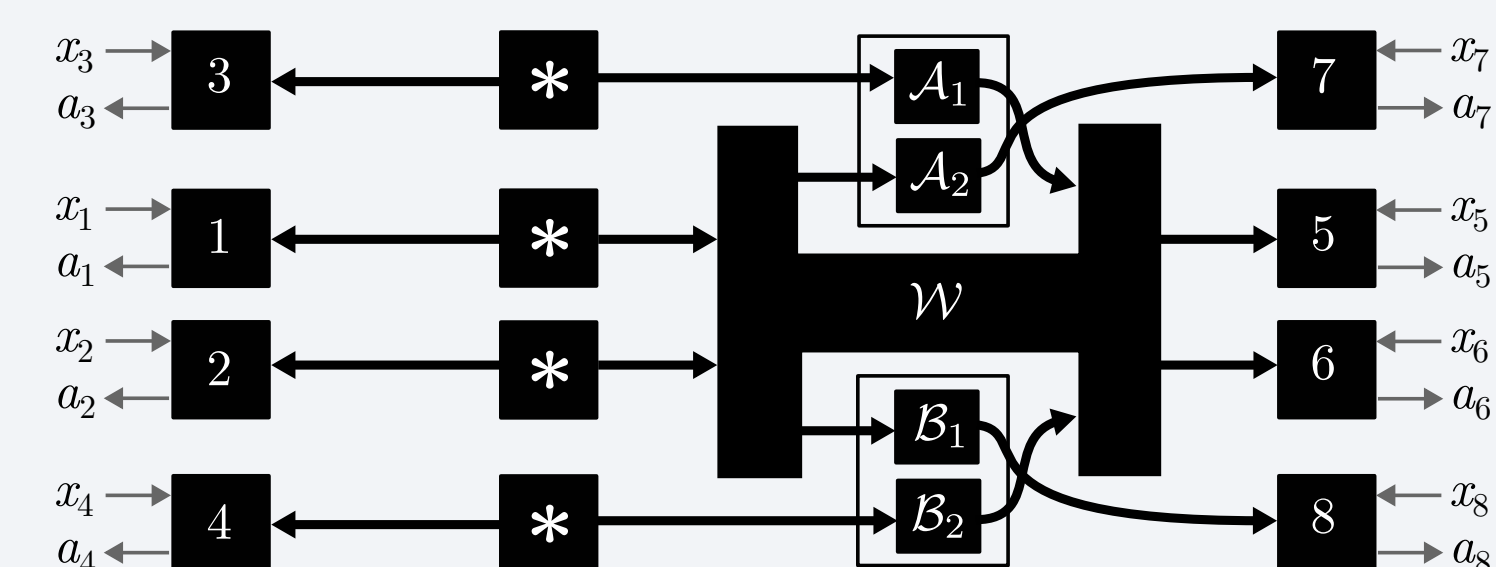
**Theorem: Self-testing the Quantum Switch**

- Given these statistics,  $\exists \Lambda_{P_c}, \Lambda_{P_t}, \Lambda_{F_c}, \Lambda_{F_t}, \Omega_A, \Omega_B$  such that:



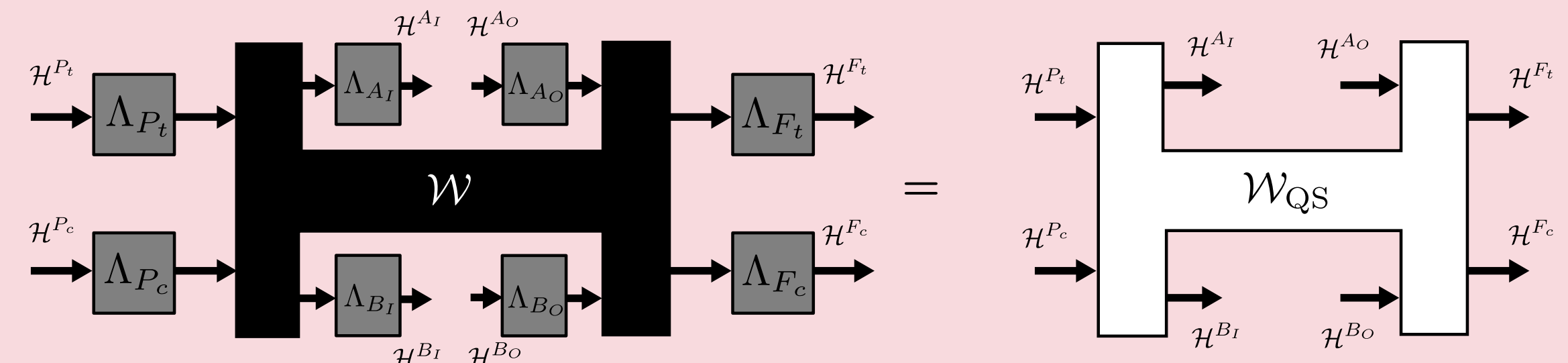
### A finer self-testing statement

- Self-test is up to **local embedding combs** at slots for  $\mathcal{A}$  and  $\mathcal{B}$ 
  - Impossible to have a finer statement if  $\mathcal{A}$  and  $\mathcal{B}$  considered whole black boxes
- **Observation:** In the black-box setting, we control network connectivity
  - Connect separate devices to  $\mathcal{H}^{A_I}, \mathcal{H}^{A_O}, \mathcal{H}^{B_I}, \mathcal{H}^{B_O}$
  - $\mathcal{A}$  and  $\mathcal{B}$  often considered closed laboratories: Alice, Bob plug devices as desired



**Theorem: Self-testing the Quantum Switch up to local channels**

- Given the same statistics,  $\exists \Lambda_{P_c}, \Lambda_{P_t}, \Lambda_{A_I}, \Lambda_{A_O}, \Lambda_{B_I}, \Lambda_{B_O}, \Lambda_{F_c}, \Lambda_{F_t}$  such that:



### Conclusions and open questions

- Can identify degrees of freedom on which the implemented supermap is precisely  $\mathcal{W}_{\text{QS}}$
- Self-testing approach applicable to general supermaps, e.g. quantum combs with 1 slot
- Which causally nonseparable quantum supermaps can be self-tested?
- Simplify setting and relax assumptions

### References and acknowledgments

- [1] P. Sekatski, J.-D. Bancal, S. Wagner, and N. Sangouard, Phys. Rev. Lett. **121**, 180505 (2018).
- [2] T. van der Lugt, J. Barrett, and G. Chiribella, Nat. Comm. **14**, 5811 (2023).
- [3] H. Dourdent, A. A. Abbott, I. Šupić, and C. Branciard, (2023), arXiv:2308.12760 [quant-ph].

This work is supported by the Agence Nationale de la Recherche projects EPiQ (ANR-22-PETQ-0007, part of Plan France 2030), and TaQC (ANR-22-CE47-0012).